

# Asymptotics for the simplest generalized Jacobi polynomials recurrence coefficients from Freud's equations: numerical explorations.

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Dedicated to L. Gatteschi on the occasion of his 70<sup>th</sup> birthday  
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**Abstract.** Generalized Jacobi polynomials are orthogonal polynomials related to a weight function which is smooth and positive on the whole interval of orthogonality up to a finite number of points, where algebraic singularities occur. The influence of these singular points on the asymptotic behaviour of the recurrence coefficients is investigated.

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## 1. Weight singularities and recurrence coefficients.

The orthonormal polynomials  $p_n(x) = \gamma_n x^n + \dots$  related to the weight  $w$  satisfy the three-terms recurrence relation

$$a_{n+1}p_{n+1}(x) = (x - b_n)p_n(x) - a_n p_{n-1}(x), \quad (1)$$

with  $a_0 p_{-1}(x) \equiv 0$ .

Let  $w(x) > 0$  hold almost everywhere on the support  $[-1, 1]$ , then one knows (since 1977) that the recurrence coefficients have limits  $a_n \rightarrow 1/2$  and  $b_n \rightarrow 0$  when  $n \rightarrow \infty$  (see for instance the survey in [Nev2]).

Features of  $w$  can somehow be “read” in the sequence of the recurrence coefficients. This goes back to Stieltjes and is current practice in solid-state physics [LaG]. See [Ap] for the case of endpoints singularities. J.P. Gaspard once showed me a paper by C.Hodges [Ho] describing the influence of a mild *interior* singularity of the form

$$\begin{aligned} w(x) &\sim w(x_0) + A(x - x_0)^\gamma, & x \rightarrow x_0, x > x_0, \\ &\sim w(x_0) + B(x_0 - x)^\gamma, & x \rightarrow x_0, x < x_0, \end{aligned} \quad (2)$$

with  $x_0 \in (-1, 1)$ ,  $0 < w(x) < \infty$  on  $(-1, 1)$  and  $\gamma > 0$  (Van Hove singularity [Ho,Mart]) as

$$a_n - 1/2 \sim \xi n^{-\gamma-1} \cos(2n\theta_0 - \eta), \quad b_n \sim 2\xi n^{-\gamma-1} \cos((2n+1)\theta_0 - \eta), \quad n \rightarrow \infty$$

with  $x_0 = \cos \theta_0$ . I gave a lengthy proof of this in [Mag1], showing how  $\xi$  and  $\eta$  are related through Toeplitz determinants to the Szegő function of  $w(\cos \theta) \sin \theta$  (the analytic function  $D(z)$  in  $|z| < 1$ , with  $D(0) > 0$ , without zero in  $|z| < 1$ , and such that the boundary values satisfy  $|D(e^{i\theta})|^2 = w(\cos \theta) \sin \theta$ ) by

$$\xi e^{i\eta} = \frac{\left(\frac{\sin \theta_0}{2}\right)^{\gamma+1} \Gamma(\gamma+1)}{2\pi w(x_0)} e^{i(\theta_0 - 2 \arg(D(e^{i\theta_0})))} [Ae^{i\pi\gamma/2} - Be^{-i\pi\gamma/2}].$$

Hodges argument is based on the continued fraction expansion of

$$f(z) = \int_{-1}^1 \frac{w(t) dt}{z-t} = \frac{\mu_0}{z-b_0 - \frac{a_1^2}{z-b_1 - \dots}}$$

when  $z = x + i\varepsilon$  for small  $\varepsilon > 0$  and  $-1 < x < 1$ , but the fastest explanation is probably related to inverse scattering techniques ([VA] p.117, [NV] and references therein): let us consider the function  $q_n$  defined by the integral  $q_n(z) = \int_{-1}^1 (z-t)^{-1} p_n(t) w(t) dt$  when  $z \notin [-1, 1]$ . Remark that  $f(z) = \sqrt{\mu_0} q_0(z)$ . From (1),  $a_{n+1} q_{n+1}(z) = (z-b_n) q_n(z) - a_n q_{n-1}(z) - \sqrt{\mu_0} \delta_{n,0}$ , (with  $a_0 q_{-1}(z) \equiv 0$ ), which can be written as

$$q_{n+1}(z) - 2z q_n(z) + q_{n-1}(z) = q_{n+1}(z) - \rho(z)^{-1} q_n(z) - \rho(z) [q_n(z) - \rho(z)^{-1} q_{n-1}(z)] = \varepsilon_n(z), \quad n = 0, 1, \dots$$

with  $\varepsilon_n(z) = (1 - 2a_{n+1}) q_{n+1}(z) - 2b_n q_n(z) + (1 - 2a_n) q_{n-1}(z) - 2\sqrt{\mu_0} \delta_{n,0}$ , and where  $\rho(z)$  is the determination of  $(z + \sqrt{z^2 - 1})/2$  such that  $|\rho(z)| > 1$  when  $z \notin [-1, 1]$ . After a simple summation,  $\rho(z)^{-N} (q_{N+1}(z) - \rho(z)^{-1} q_N(z)) - \rho(z) q_0(z) = \sum_0^N \rho(z)^{-n} \varepsilon_n(z)$ . While  $z \notin [-1, 1]$  and when  $N \rightarrow \infty$ ,  $\rho(z)^{-N}$  and  $q_N(z) \rightarrow 0$ , so

$$q_0(z) = \frac{f(z)}{\sqrt{\mu_0}} = - \sum_0^\infty \frac{\varepsilon_n(z)}{\rho(z)^{n+1}} = \frac{2\sqrt{\mu_0}}{\rho(z)} + \sum_0^\infty \frac{(2a_n - 1)(q_{n-1}(z) + \rho(z) q_n(z)) + 2b_n q_n(z)}{\rho(z)^{n+1}}$$

From the known asymptotic behaviour  $q_n(z) \sim \frac{2(2\pi)^{1/2} D(\rho(z)^{-1})}{(\rho(z) - \rho(z)^{-1}) \rho(z)^n}$  for large  $n$  [Ba], we have a first approximate expansion of  $(2\pi)^{-1/2} D(\rho(z)^{-1}) (\rho(z) - \rho(z)^{-1}) q_0(z)$  as a series of negative powers of  $\rho(z)$ :

$$(2\pi)^{-1/2} D(\rho(z)^{-1}) (\rho(z) - \rho(z)^{-1}) q_0(z) \approx \sum_0^\infty \frac{2(2a_n - 1)}{\rho(z)^{2n}} + \frac{2b_n}{\rho(z)^{2n+1}}$$

Under sufficiently strong conditions, this remains valid when  $z = x + i\varepsilon$ ,  $\varepsilon \rightarrow 0$ ,  $\varepsilon > 0$ ,  $-1 \leq x = \cos(\theta) \leq 1$ ,  $\rho(z) \rightarrow e^{i\theta}$  (with  $0 \leq \theta \leq \pi$ ), so that  $2a_n - 1$  and  $b_n$  behave like the *Fourier coefficients* of order  $2n$  and  $2n+1$  of a function whose singularities on  $0 < \theta < \pi$  are related to the singularities of  $w(\cos \theta) \sin \theta$ . Singularities of type  $|\theta - \theta_0|^\gamma$  correspond indeed to  $n^{-1-\gamma} \cos(n\theta_0 + \text{const.})$  behaviour in the  $n^{\text{th}}$  Fourier coefficient.

This analysis is no more valid for stronger singularity  $w(x_0) = 0$  or  $\infty$  because the approximations done near  $x_0$  are no more valid, subtle important effects take place in neighbourhoods of length about  $1/n$  of the singular points, see the famous [NevIII].

It is therefore not useless to have a close look at the simplest orthogonal polynomials related to weights with interior singularities. Special singular positions have been worked (in sieved polynomials theory etc.), but here is something related to an *arbitrary* position:

## 2. Freud's equations for the simplest generalized Jacobi polynomials.

Let

$$\begin{aligned} w(x) &= B(1-x)^\alpha(x-x_0)^\gamma(1+x)^\beta \quad x \in [-1, x_0], \\ &= A(1-x)^\alpha(x-x_0)^\gamma(1+x)^\beta \quad x \in [x_0, 1], \end{aligned} \quad (3)$$

with  $-1 < x_0 < 1$ ,  $A$  and  $B > 0$ ,  $\alpha, \beta$  and  $\gamma > -1$ .

The case  $A = B, \alpha = \beta, x_0 = 0$  has the simple solution deduced from Jacobi polynomials  $a_n^2 = (n + 2\alpha + \gamma \text{ odd}(n))(n + \gamma \text{ odd}(n))/[(2n + 2\alpha + \gamma + 1)(2n + 2\alpha + \gamma - 1)]$ ,  $b_n = 0$ , where  $\text{odd}(n) = (1 - (-1)^n)/2$ . When  $n \rightarrow \infty$ , this case shows the asymptotic behaviour  $a_n \sim \frac{1}{2} - \frac{(-1)^n \gamma}{4n}$ ,  $b_n = 0$ .

The  $O(1/n)$  term is definitely related to the  $|x|^\gamma$  behaviour of the weight near 0, as shown by P. Nevai (Theorem 4 of Section 7 of [Nev]): if  $w$  is even on  $[-1, 1]$ , with  $w(x)|x|^{-\gamma}$  positive and continuously derivable on  $(-1, 1)$ , then  $a_n = 1/2 - \gamma(-1)^n/(4n) + o(1/n)$  when  $n \rightarrow \infty$ .

We now try to investigate the recurrence coefficients  $a_n, b_n$  when the weight is (3). This weight is a *semi-classical* weight, as  $w'/w$  is the same rational function almost everywhere on the support  $[-1, 1]$  of  $w$ .

Semi-classical orthogonal polynomials have a rich differential structure, according to a theory going as far as Laguerre [BeR, GaN, Lag, Mag2, Mag3, Sho].

Freud [Fr] showed how to deduce recurrence coefficients asymptotics from special identities. For a general semi-classical weight satisfying  $w'(x)/w(x) = 2V(x)/W(x)$  with  $W(x)w(x) \rightarrow 0$  when  $x$  tends to any endpoint of the support  $S$  of  $w$  (Shohat's conditions [Sho]), we find these identities (*Freud's equations*) by expanding

$$\begin{aligned} 0 &= \int_S [W(x)w(x)p_n(x)p_{n-k}(x)]' dx \\ &= \int_S W(x)w(x)p_n'(x)p_{n-k}(x) dx + \int_S W(x)w(x)p_n(x)p_{n-k}'(x) dx + \\ &\quad + \int_S W'(x)w(x)p_n(x)p_{n-k}(x) dx + \int_S 2V(x)w(x)p_n(x)p_{n-k}(x) dx \end{aligned}$$

for  $k = 0, 1$ , remarking that any integral  $\int_S P(x)w(x)p_n(x)p_{n-k}(x)dx$  where  $P$  is a polynomial, is an expression involving  $a_n, b_n, a_{n\pm 1}, b_{n\pm 1}$ , etc. according to  $k$  and the degree of  $P$ , and using  $p_n' = np_{n-1}/a_n + (b_0 + \dots + b_{n-1} - nb_{n-1})p_{n-2}/(a_{n-1}a_n) + \dots$  (see [BeR]).

Reduction to even measure: one considers the orthonormal polynomials  $\{\tilde{p}_n\}$  with respect to the even weight

$$\tilde{w}(x) = 2|x|w(2x^2 - 1), \quad \text{for } -1 < x < 1.$$

Then,  $\tilde{p}_{2n}(x) = p_n(2x^2 - 1)$ , and one recovers the recurrence relation for the  $p_n$ 's by contracting the recurrence relation for the  $\tilde{p}_n$ 's:

$$\begin{aligned} \tilde{a}_{n+1}\tilde{p}_{n+1}(x) &= x\tilde{p}_n(x) - \tilde{a}_n\tilde{p}_{n-1}(x) \quad \Rightarrow \\ \tilde{a}_{2n+1}\tilde{a}_{2n+2}\tilde{p}_{2n+2}(x) &= (x^2 - \tilde{a}_{2n}^2 - \tilde{a}_{2n+1}^2)\tilde{p}_{2n}(x) - \tilde{a}_{2n-1}\tilde{a}_{2n}\tilde{p}_{2n-2}(x), \end{aligned}$$

so:

$$a_n = 2\tilde{a}_{2n-1}\tilde{a}_{2n}, \quad b_n = -1 + 2\tilde{a}_{2n}^2 + 2\tilde{a}_{2n+1}^2 \quad (4)$$

This allows to work with the single sequence  $\{\tilde{a}_n\}$  instead of the two sequences  $\{a_n\}, \{b_n\}$ . Here,

$$\begin{aligned} \tilde{w}(x) &= 2|x|w(2x^2 - 1) = \tilde{B}|x|^{2\beta+1}(\tilde{x}_0^2 - x^2)^\gamma(1 - x^2)^\alpha & \text{for } |x| < |\tilde{x}_0|, \\ &= \tilde{A}|x|^{2\beta+1}(x^2 - \tilde{x}_0^2)^\gamma(1 - x^2)^\alpha & \text{for } |\tilde{x}_0| < |x| < 1, \end{aligned}$$

where  $\tilde{x}_0$  is the positive root of  $2\tilde{x}_0^2 - 1 = x_0$ ,  $\tilde{A} = 2^{\alpha+\beta+\gamma+1}A$ ,  $\tilde{B} = 2^{\alpha+\beta+\gamma+1}B$ .

So,  $\tilde{W}(x) = x(x^2 - \tilde{x}_0^2)(x^2 - 1) = x^5 - (\tilde{x}_0^2 + 1)x^3 + \tilde{x}_0^2x$  and  $2\tilde{V}(x) = (2\alpha + 2\beta + 2\gamma + 1)x^4 - [2\alpha\tilde{x}_0^2 + (2\beta + 1)(\tilde{x}_0^2 + 1) + 2\gamma]x^2 + (2\beta + 1)\tilde{x}_0^2$ .

The equations for the  $\tilde{a}_n$ 's now follow from Freud's method for even weights, expanding  $\tilde{W}\tilde{w}' = 2\tilde{V}\tilde{w}$  as

$$\begin{aligned} \int_S \frac{\tilde{W}(x)}{x} \tilde{w}(x) \tilde{p}'_n(x) \tilde{p}_{n-1}(x) + \int_S \frac{\tilde{W}(x)}{x} \tilde{w}(x) \tilde{p}_n(x) \tilde{p}'_{n-1}(x) + \int_S \left( \frac{\tilde{W}(x)}{x} \right)' \tilde{w}(x) \tilde{p}_n(x) \tilde{p}_{n-1}(x) + \\ + \int_S \frac{2\tilde{V}(x)}{x} \tilde{w}(x) \tilde{p}_n(x) \tilde{p}_{n-1}(x) = 0, \end{aligned}$$

using  $\int_S \tilde{w}(x) \tilde{p}_i(x) \tilde{p}_j(x) dx = \delta_{i,j}$ , the recurrence relations (1) giving  $\int_S x \tilde{w}(x) \tilde{p}_n(x) \tilde{p}_{n-1}(x) dx = \tilde{a}_n$ ,  $\int_S x^2 \tilde{w}(x) (\tilde{p}_n(x))^2 dx = \tilde{a}_n^2 + \tilde{a}_{n+1}^2$ ,  $\int_S \tilde{w}(x) x^{-1} \tilde{p}_n(x) \tilde{p}_{n-1}(x) dx = \text{odd}(n)/\tilde{a}_n$ , etc., and

$$\begin{aligned} \tilde{p}'_n &= \frac{n}{\tilde{a}_n} \tilde{p}_{n-1} + \frac{2 \sum_1^{n-1} \tilde{a}_k^2 - n \tilde{a}_{n-1}^2}{\tilde{a}_{n-2} \tilde{a}_{n-1} \tilde{a}_n} \tilde{p}_{n-3} + \\ &+ \frac{n \tilde{a}_{n-3}^2 \tilde{a}_{n-1}^2 - 2(\tilde{a}_{n-3}^2 + \tilde{a}_{n-2}^2 + \tilde{a}_{n-1}^2) \sum_1^{n-1} \tilde{a}_k^2 + 2 \sum_1^{n-1} (\tilde{a}_k^4 + 2\tilde{a}_k^2 \tilde{a}_{k-1}^2)}{\tilde{a}_{n-4} \tilde{a}_{n-3} \tilde{a}_{n-2} \tilde{a}_{n-1} \tilde{a}_n} \tilde{p}_{n-5} + \dots, \end{aligned}$$

one finally finds

$$\begin{aligned} 2(n + \alpha + \beta + \gamma + 2)\tilde{a}_n^2(\tilde{a}_{n-1}^2 + \tilde{a}_n^2 + \tilde{a}_{n+1}^2) - 2[\alpha\tilde{x}_0^2 + (n + \beta + 1)(\tilde{x}_0^2 + 1) + \gamma]\tilde{a}_n^2 + \\ + 2(2\tilde{a}_n^2 - \tilde{x}_0^2 - 1) \sum_{j=1}^{n-1} \tilde{a}_j^2 + n\tilde{x}_0^2 - 2\tilde{a}_n^2 \tilde{a}_{n-1}^2 + 2 \sum_{j=1}^{n-1} (\tilde{a}_j^4 + 2\tilde{a}_j^2 \tilde{a}_{j-1}^2) + (2\beta + 1)\tilde{x}_0^2 \text{ odd}(n) = 0, \\ n = 1, 2, \dots \end{aligned} \quad (5)$$

( $\tilde{a}_0 = 0$ ).

We see how any  $\tilde{a}_n$  can be computed from the value of  $\tilde{a}_1$ , which is the degree of freedom reflecting that the same equations (5) hold for any choice of  $\tilde{A}$  and  $\tilde{B}$  in the weight  $\tilde{w}$ . Actually  $\tilde{a}_1$  is linked to the ratio  $\tilde{A}/\tilde{B}$  by

$$\tilde{a}_1^2 = \frac{\tilde{\mu}_2}{\tilde{\mu}_0} = \frac{\tilde{B} \int_{|x| < \tilde{x}_0} |x|^{2\beta+3} (\tilde{x}_0^2 - x^2)^\gamma (1 - x^2)^\alpha dx + \tilde{A} \int_{|x| > \tilde{x}_0} |x|^{2\beta+3} (x^2 - \tilde{x}_0^2)^\gamma (1 - x^2)^\alpha dx}{\tilde{B} \int_{|x| < \tilde{x}_0} |x|^{2\beta+1} (\tilde{x}_0^2 - x^2)^\gamma (1 - x^2)^\alpha dx + \tilde{A} \int_{|x| > \tilde{x}_0} |x|^{2\beta+1} (x^2 - \tilde{x}_0^2)^\gamma (1 - x^2)^\alpha dx} \quad (6)$$

Numerical experiments show that  $\tilde{a}_2, \tilde{a}_3, \dots$  can be computed in a stable way from  $\tilde{a}_1$  simply by considering (5) as an equation for  $\tilde{a}_{n+1}$  when  $\tilde{a}_1, \dots, \tilde{a}_n$  are known:

$$\begin{aligned} \tilde{a}_{n+1}^2 = & \frac{\alpha\tilde{x}_0^2 + (n + \beta + 1)(\tilde{x}_0^2 + 1) + \gamma}{N} - 2\frac{\sum_1^{n-1} \tilde{a}_k^2}{N} + \\ & + \frac{2(\tilde{x}_0^2 + 1) \sum_1^{n-1} \tilde{a}_k^2 - n\tilde{x}_0^2 - 2 \sum_1^{n-1} (\tilde{a}_k^4 + 2\tilde{a}_k^2 \tilde{a}_{k-1}^2) - (2\beta + 1)\tilde{x}_0^2 \text{ odd}(n)}{2N\tilde{a}_n^2} + \frac{\tilde{a}_{n-1}^2}{N} - \tilde{a}_n^2 - \tilde{a}_{n-1}^2 \end{aligned} \quad (7)$$

with  $N = n + \alpha + \beta + \gamma + 2$ ,  $n = 1, 2, \dots$

### 3. Asymptotic estimates.

Some people call these tricks “special refinements”; others call them “kludges”.

D.E. Knuth

Putting an almost constant  $\tilde{a}_n^2 \approx \tilde{a}^2$  in (5) gives two possible asymptotic matches  $\tilde{a}^2 = 1/4$  and  $\tilde{a}^2 = \tilde{x}_0^2/4$ , corresponding to weights with support  $[-1, 1]$  and  $[-\tilde{x}_0, \tilde{x}_0]$  (the latter when  $A = 0$ ). More complicated behaviours are expected to hold when  $\tilde{x}_0$  is complex [GaN] and one should be able to establish correct asymptotic behaviours from (5) alone, but this has not yet been achieved (the answer to [GaN] was given in [N] with other techniques). Even when one knows that  $\tilde{a}_n^2 \rightarrow 1/4$ , some amount of guesswork will still be needed. Let  $\tilde{a}_n^2 = \frac{1}{4} + y_n$ . We know from Szegő's theory ([Sz] chap. 12) that  $\sum_1^{n-1} \tilde{a}_k^2 = \frac{n}{4} + \xi + z_n$  and  $\sum_1^{n-1} (\tilde{a}_k^4 + 2\tilde{a}_k^2 \tilde{a}_{k-1}^2) = \frac{3n}{16} + \eta + u_n$  with  $z_n$  and  $u_n \rightarrow 0$  (See also [Nev2] p.91). So,

$$\begin{aligned} y_{n+1} = & -\frac{1}{4} + \frac{\alpha\tilde{x}_0^2 + (n + \beta + 1)(\tilde{x}_0^2 + 1) + \gamma}{N} - 2\frac{\frac{n}{4} + \xi + z_n}{N} + \\ & + \frac{2(\tilde{x}_0^2 + 1)(\frac{n}{4} + \xi + z_n) - n\tilde{x}_0^2 - 2(\frac{3n}{16} + \eta + u_n) - (2\beta + 1)\tilde{x}_0^2 \text{ odd}(n)}{2N(\frac{1}{4} + y_n)} + \frac{\frac{1}{4} + y_{n-1}}{N} - \frac{1}{2}y_n - y_{n-1}. \end{aligned}$$

We compute now  $\xi$  and  $\eta$ : from the Szegő theory, let  $\tilde{\phi}_n(z) = \tilde{\kappa}_n z^n + \tilde{\kappa}'_n z^{n-2} + \tilde{\kappa}''_n z^{n-4} + \dots$  be the orthonormal polynomials on the unit circle with respect to  $|\sin \theta| \tilde{w}(\cos \theta) = \tilde{C}(\theta) |\cos \theta|^{2\beta+1} |\cos^2 \theta - \cos^2(\theta_0/2)|^\gamma |\sin \theta|^{2\alpha+1}$ , with  $\tilde{C}(\theta) = \tilde{A}$  on  $|\theta| < \theta_0/2$  and  $|\theta - \pi| < \theta_0/2$  and  $\tilde{C}(\theta) = \tilde{B}$  elsewhere. The Szegő function  $\tilde{D}(z)$  whose boundary values must be  $|\tilde{D}(e^{i\theta})| = \sqrt{\tilde{w}(\cos \theta) |\sin \theta|}$  is found by inspection to be

$$\tilde{D}(z) = \tilde{\kappa}^{-1} (1 - z^2)^{\alpha+1/2} (1 + z^2)^{\beta+1/2} (1 - e^{-i\theta_0} z^2)^{\gamma/2+i\lambda} (1 - e^{i\theta_0} z^2)^{\gamma/2-i\lambda},$$

with  $\tilde{\kappa} = 2^{\alpha+\beta+\gamma+1} \tilde{B}^{(\theta_0-\pi)/(2\pi)} \tilde{A}^{-\theta_0/(2\pi)}$  and  $\lambda = (2\pi)^{-1} \log(\tilde{B}/\tilde{A})$ . We find the limit values of  $\tilde{\kappa}'_n$  and  $\tilde{\kappa}''_n$  from the expansion of  $1/\tilde{D}$ :

$$\begin{aligned} \tilde{\kappa}'_n / \tilde{\kappa}_n & \rightarrow \alpha - \beta + \gamma x_0 + 2\lambda \sin \theta_0, \\ \tilde{\kappa}''_n / \tilde{\kappa}_n & \rightarrow (\alpha - \beta)^2 / 2 + (\alpha + \beta + 1) / 2 + (\alpha - \beta)(\gamma x_0 + 2\lambda \sin \theta_0) + \\ & + (\gamma(\gamma + 2) / 4 - \lambda^2) \cos(2\theta_0) + \lambda(\gamma + 1) \sin(2\theta_0) + \gamma^2 / 4 + \lambda^2, \end{aligned}$$

used in

$$\begin{aligned}
\tilde{P}_n(x) &= \frac{\tilde{p}_n(x)}{\tilde{\gamma}_n} = \\
&= x^n - \left( \sum_1^{n-1} a_k^2 \right) x^{n-2} + \left[ \left( \sum_{k=1}^{n-1} \tilde{a}_k^2 \right)^2 - \left( \sum_{k=1}^{n-1} \tilde{a}_k^4 \right) - 2 \left( \sum_{k=1}^{n-2} \tilde{a}_k^2 \tilde{a}_{k+1}^2 \right) \right] x^{n-4}/2 + \dots \\
&= \frac{z^{-n} \tilde{\phi}_{2n}(z) + z^n \tilde{\phi}_{2n}(z^{-1})}{2^n (\tilde{\kappa}_{2n} + \tilde{\phi}_{2n}(0))} \\
&\sim \frac{T_n(x)}{2^{n-1}} + \frac{\tilde{\kappa}' T_{n-2}(x)}{\tilde{\kappa} 2^{n-1}} + \frac{\tilde{\kappa}'' T_{n-4}(x)}{\tilde{\kappa} 2^{n-1}} + \dots \\
&\sim x^n - \frac{n - \tilde{\kappa}'/\tilde{\kappa}}{4} x^{n-2} + \frac{n(n-3)/2 - (n-2)\tilde{\kappa}'/\kappa + \tilde{\kappa}''/\tilde{\kappa}}{16} x^{n-4} + \dots
\end{aligned}$$

$$(z + z^{-1} = 2x).$$

So,

$$\xi = \lim_{n \rightarrow \infty} \sum_1^{n-1} \tilde{a}_k^2 - n/4 = -\kappa'/(4\kappa) = -(\alpha - \beta + \gamma x_0 + 2\lambda \sin \theta_0)/4, \quad (8)$$

$$\eta = \lim_{n \rightarrow \infty} \sum_1^{n-1} (\tilde{a}_k^4 + 2\tilde{a}_k^2 \tilde{a}_{k-1}^2) - 3n/16 = ((\kappa'/\kappa)^2 - 4\kappa'/\kappa - 2\kappa''/\kappa)/16.$$

The equation for the  $y_n$ 's reduces to

$$\begin{aligned}
y_{n+1} - 2x_0 y_n + y_{n-1} &= \frac{(x_0 + 1)(\beta + 1/2)(-1)^n + 2(x_0 + 2)z_n + y_{n-1} - 4u_n}{N} - \\
&- \frac{16y_n}{1 + 4y_n} \frac{(2x_0 + 1)/4 - \lambda \sin \theta_0 + (x_0 + 1)(\beta + 1/2)(-1)^n + 2(x_0 + 3)z_n - 4u_n}{4N} \\
&- \frac{4(2x_0 + 1)y_n^2}{1 + 4y_n} \quad (9)
\end{aligned}$$

where  $z_n = -\sum_n^\infty y_k$  and  $u_n = -y_{n-1}/2 - \sum_n^\infty (3y_k/2 + y_k^2 + 2y_k y_{k-1})$ .

The form (9) should give hints on the behaviour of  $y_n$  when  $n \rightarrow \infty$ . No proof will be attempted here, only reasonable asymptotic matching and numerical checks. Use of Painlevé-like differential equations in  $\theta_0$  is another method of investigation which could be used in the future (see [Mag3]).

The right-hand side is small, even with respect to the  $y$ 's, so  $y_{n+1} - 2x_0 y_n + y_{n-1}$  is small, and this suggests a  $\exp(\pm in\theta_0)$  behaviour somewhere. However, I still don't have a tight proof that  $y_n$ ,  $z_n$  and  $u_n$  are  $O(1/n)$ . Assuming  $y_n = K_1(-1)^n/n + K_2 e^{in\theta_0}/n^{\zeta_2} + K_3 e^{-in\theta_0}/n^{\zeta_3}$ , matching the two sides gives  $K_1 = -\beta/2 - 1/4$ ,  $\zeta_{2,3} = 1 \pm 2i\lambda$ . Numerical checks have been performed on the form

$$\tilde{a}_n^2 - 1/4 = y_n = -(\beta + 1/2)(-1)^n/(2n) + K \cos(n\theta_0 - 2\lambda \log n - \varphi)/n + o(1/n) \quad (10)$$

where  $K$  and  $\varphi$  are unknown functions of  $\alpha$ ,  $\beta$ ,  $\gamma$ ,  $\lambda$  and  $x_0$ . Given  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $x_0$ , the algorithm first performs (7) with several trial starting values  $\tilde{a}_1^2$  and computes the corresponding  $\lambda$  from (8), allowing the determination of the coefficients in (6). It is then possible to run (7) for a requested value of  $\lambda$ , and to estimate  $K$  and  $\varphi$  in (10) from numerical values of  $\tilde{a}_n^2$  for large  $n$  (up to the 10000-100000 range). Very satisfactory empirical formulas

for  $K$  and  $\varphi$  appear to be  $K = (\gamma^2/4 + \lambda^2)^{1/2} \sin(\theta_0/2)$  and  $\varphi = (\alpha + 1 + \gamma/2)\pi - (\alpha + \beta + \gamma + 1/2)\theta_0 + 2\lambda \log(2 \sin \theta_0) - 2 \arg \Gamma(\gamma/2 + i\lambda) - \arg(\gamma/2 + i\lambda)$ .

Whence, from (4):  $a_n - 1/2 \sim y_{2n-1} + y_{2n}$ ,  $b_n \sim 2(y_{2n} + y_{2n+1})$ , the

**Conjecture.** The recurrence coefficients related to the simplest generalized Jacobi weight (3) satisfy

$$a_n = \frac{1}{2} - \frac{M}{n} \cos [2n\theta_0 - 2\lambda \log(4n \sin \theta_0) - \Phi] + o\left(\frac{1}{n}\right),$$

$$b_n = -\frac{2M}{n} \cos [(2n+1)\theta_0 - 2\lambda \log(4n \sin \theta_0) - \Phi] + o\left(\frac{1}{n}\right),$$

when  $n \rightarrow \infty$ , where  $x_0 = \cos \theta_0$ ,  $0 < \theta_0 < \pi$ ,  $\lambda = \log(B/A)/(2\pi)$ ,  $M = \frac{1}{2}(\gamma^2/4 + \lambda^2)^{1/2} \sin \theta_0$ ,  $\Phi = (\alpha + \gamma/2)\pi - (\alpha + \beta + \gamma)\theta_0 - 2 \arg \Gamma(\gamma/2 + i\lambda) - \arg(\gamma/2 + i\lambda)$ .

Here is a sample of the numerical check:  $K$  and  $\varphi$  are extracted from the form (10) on a sample of  $\tilde{a}_n^2$ , with  $n$  going up to 500000. The values of  $K$  and  $\lambda$  (from (8)) are quite stable, but things are not so easy with  $\varphi$  (the “phase” column). Finally, the line “check” contains the computed values of  $K$  and  $\varphi$ , and a new check of  $\lambda$  through Turán determinant weight function reconstruction yielding  $\tilde{A}$  and  $\tilde{B}$ .

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Script started on Wed Mar 2 17:38:49 1994

/u18/grpanma/magnus/jacobi@ux12 [1] #ls -l

total 1224

```

-rwxr-xr-x  1 magnus      473298 Mar  2 17:30 a.out
-rw-r--r--  1 magnus        2226 Feb 16 15:24 alnrel.f
-rwxr-xr-x  1 magnus         18 Jan  3 14:15 batchsub
-rw-r----- 1 magnus       1183 Feb 16 08:58 c8lgmc.fz
-rw-r----- 1 magnus       2087 Feb 16 08:58 c9lgmc.f
-rw-r--r--  1 magnus        259 Feb 16 15:25 carg.f
-rw-r----- 1 magnus      21751 Feb 16 16:41 clngam.f
-rw-r--r--  1 magnus     14412 Feb 16 16:42 clngam.o
-rw-r----- 1 magnus       2645 Feb 16 15:39 clngamold.f
-rw-r----- 1 magnus        975 Feb 16 08:58 clnrel.f

```



```

-rw-r--r-- 1 magnus      141 Feb 16 16:35 cmlxt.f
-rw-r--r-- 1 magnus     1061 Feb 16 15:52 csevl.f
-rw-r----- 1 magnus      471 Feb 16 08:58 entsrc.fz
-rw-r--r-- 1 magnus     7679 Mar  2 17:29 gjacobi.f
-rw-r--r-- 1 magnus    14554 Mar  2 17:30 gjacobi.o
-rw-r--r-- 1 magnus     6394 Feb  9 09:07 gjacobi4.f
-rw-r----- 1 magnus     4153 Feb 16 08:59 indexfn.z
-rw-r--r-- 1 magnus      861 Feb 16 15:52 inits.f
-rw-r--r-- 1 magnus      167 Feb 16 16:46 lgam.f
-rw-r--r-- 1 magnus     1325 Feb 16 16:46 lgam.o
-rwxr-xr-x 1 magnus    452258 Mar  2 13:50 lgam.out
-rw-r----- 1 magnus    10873 Feb 16 08:58 macharb.f
-rw-r----- 1 magnus     4759 Feb 16 08:58 r1macha.fz
-rw-r----- 1 magnus     8318 Feb 16 08:58 r1machb.f
-rw-r----- 1 magnus     2800 Feb 16 08:58 seterr.f
-rw-r----- 1 magnus      319 Feb 16 08:58 seteru.f

```

```
/u18/grpanma/magnus/jacobi@ux12 [2] #cat gjacobi.f
```

```

c  gjacobi.for
c      program gjacobi
c  simplest generalized Jacobi
c
c      w(x) = B (1-x)^alpha (1+x)^beta (x0-x)^gamma  -1<x<x0
c             = A (1-x)^alpha (1+x)^beta (x-x0)^gamma  x0<x<1
c
c  -> w tilde(y)= B tilde |y|^(2beta+1) (y0-y)^gamma (1-y^2)^alpha
c                                     -y0<|y|<y0
c  -> w tilde(y)= A tilde |y|^(2beta+1) (y-y0)^gamma (1-y^2)^alpha
c                                     y0<|y|<1
c
c  x=2y^2-1
c
c  the program produces squares of recurrence coefficients atilde n
c
c  of   atilde n+1 ptilde n+1 (y) = y ptilde n (y) - atilde n ptilde n-1 (y)
c
c  where the ptilde s are the orthonormal polynomials related to w tilde
c
c  input: first, period of display of atilde n, number of values of y
c          where to compute polynomials, and interval (this latter part
c          is alpha test, enter 0 0 0)
c
c  then: alpha, beta, gamma, y0 , n max
c
c  then, after preliminary calculations have been done,
c          lambda = log(B/A) / (2 pi)

```

```

c
integer oddn,n,nmax,i,signg,recomp,period,npol
double precision al,be,ga,y0,atn,atnp1,pi,aux,shalf,cth,sth,
&      at2,at2p1,at2m1,one,half,
&      th0, x, fn, cnt,snt,s1,s2,snp1,cnp1,xx,xy,yy,lam1,lam2
&      ,aux1,at,bt,ba(4),at21(4),c1,c2,c3,thx0,lam0,phi0,phi
&      ,phi1,lsth
c or, if you can afford it (useful if |lambda|>5) :
c real*16 al,be,ga,y0,atn,atnp1,pi,aux,shalf,cth,sth,
c &      at2,at2p1,at2m1,one,half,
c &      th0, x, fn, cnt,snt,s1,s2,snp1,cnp1,xx,xy,yy,lam1,lam2
c &      ,aux1,at,bt,ba(4),at21(4),c1,c2,c3,thx0,lam0,phi0,phi
c &      ,phi1,lsth
double precision xv(502),polv(502,3),dtur
c complex gamma function from W. Fullerton's fnlib
c available by anonymous ftp on netlib.att.com cd/netlib/fn
complex clngam,z
one=1
half=one/2
pi=4*atan(one)
print *,' period of display of coeff., number of abs., extr. abs.? '
read *,period,npol,xv1,xv2
c recurrence coefficients will be printed with step period,
c polynomials and weight will be computed at xv1,...,xv2
npol=min(npol,500)
1 print *,' input alpha,beta,gamma,xtilde0,nmax (stop if nmax<=0)'
read *,al,be,ga,y0,nmax
if(nmax.le.0)stop
print '(4f14.8,i8,f9.4)',al,be,ga,y0,nmax
signg=1
if(ga.lt.0.0)signg=-1
shalf=sqrt(1-y0*y0)
cth=2*y0*y0-1
sth=2*shalf*y0
th0=acos(cth)
phi0=pi*(2*al+2+ga)/2 - (al+be+ga+half)*th0
lsth=log(2*sth)
xv(1)=y0/2
xv(2)=(1+y0)/2
do 100 i=1,npol
100 xv(i+2)=xv1+(i-1)*(xv2-xv1)/(npol-0.999)
itba=1
at2=y0*y0
at21(1)=at2
c 3 satisfactory trial values of atilde square(1) are needed to
c establish the homographic relation atilde square(1)=c1+c2B/(A+c3B)

```

```

11  if(itba.eq.1)goto 12
    if(itba.le.3)then
      if(at21(itba).le.0)at21(itba)=0.0001
      if(at21(itba).ge.1)at21(itba)=0.9999
      at21(itba)=0.99*at21(itba)+0.01*y0*y0
c to be tried until all the atilde square(n) are positive
      at2=at21(itba)
    else
      print *, ' ? lambda ( stop if > 999)'
      read *, lam0
      if(lam0.gt.999)goto 1
      ba(4)=exp(2*pi*lam0)
      at2=c1+c2/(c3+1/ba(4))
      at21(4)=at2
    endif
12  cnt=cth
    snt=sth
    s1=0
    s2=0
    at2m1=0
    oddn=1
    if(itba.ge.4)
&print *, '  n      atilde2(n)      K      phase      lambda      lambda'
      if(itba.ge.4)print '(i8,f14.10)',1,at2
      atn=sqrt(at2)
      do 13 i=1,npol+2
        polv(i,1)=1
13  polv(i,2)=xv(i)/atn
c main engine for recurrence coefficients
      do 10 n=1,nmax
        aux=al*y0*y0+(n+be+1)*(y0*y0+1)+ga+( (y0*y0+1)/at2 -2 )*s1
&      -(n*y0*y0)/(2*at2)+at2m1-s2/at2 -(be+half)*y0*y0*oddn/at2
        at2p1=aux/(n+al+be+ga+2) -at2m1-at2
c updating the two sums
        s1=s1+at2
        s2=s2+at2*( at2+2*at2m1 )
        snp1=snt*cth+cnt*sth
        cnp1=cnt*cth-snt*sth
c check asymptotic formula when needed to display
        if((n.lt.2).or.(at2p1.lt.0).or.(mod(n+1,period).eq.0).or.
&      (n.eq.nmax)) then
          aux=(n+1)*(at2p1-half/2)+(oddn-half)*(be+half)
c aux should be about K cos( (n+1)theta0 -phi )
c with phi slowly varying with n
          aux1=n*(at2-half/2)-(oddn-half)*(be+half)
c K cos(phi):

```

```

      xx=(aux1*snp1-aux*snt)/sth
c  K sin(phi):
      yy=(aux*cnt-aux1*cnp1)/sth
c phi (between -pi and pi )
      phi=atan2(yy,xx)
c K
      xx=sqrt(xx*xx+yy*yy)
c 2 ways to estimate lambda
      lam1=(-4*s1+n+1-al+be-ga*cth)/(2*sth)
      lam2=half*(-16*s2+3*(n+1)-5*al+3*be-1-ga*(2*cth*cth+4*cth-1))/
&      (2*sth*(cth+2))
      fn=n+1
c phi -2 lambda log n
      fn=phi-2*lam1*log(fn)
      fn=fn/(2*pi)
      ifn=fn
      fn=fn-ifn
      if(fn.lt.-half)fn=fn+1
      if(fn.gt. half)fn=fn-1
      phi=2*pi*fn
c Theta n(x tilde 0):
c thx0=4*s1+2*(n+al+be+ga)*at2+2*(n+2+al+be+ga)*at2p1
c &      -2*(n+be+ga+1-ga*y0*y0)
c thx0=thx0*y0
      if(itba.ge.4)
&print '(i8,f14.10,4f10.5)',n+1,at2p1,xx,phi,lam1,lam2
      endif
      if(at2p1.le.0)then
          if(itba.ge.4)
&      print *,' atildesquare <0, try a less extreme lambda'
          goto 11
      endif
c values of orthogonal polynomial of degree n+1
      atnp1=sqrt(at2p1)
      do 25 i=1,npol+2
25 polv(i,3)=( xv(i)*polv(i,2)-atn*polv(i,1) )/atnp1
c prepare next step
      if(n.lt.nmax)then
          snt=snp1
          cnt=cnp1
          oddn=1-oddn
          at2m1=at2
          at2=at2p1
          atn=atnp1
          do 251 i=1,npol+2
          polv(i,1)=polv(i,2)

```

```

251 polv(i,2)=polv(i,3)
    endif
10 continue
    if(itba.ge.4)then
c check : weight function reconstruction
c w(x)= 2 sqrt(1-x^2) /(pi * dtur)
    do 20 i=1,npol+2
        x=xv(i)
        dtur=polv(i,2)**2-atnp1*polv(i,1)*polv(i,3)/atn
        fn=1/(dtur*pi* x**(2*be+1) *(abs(y0*y0-x*x))**ga
&          *(1-x*x)**(al-half))
        if(i.eq.1)bt=fn
        if(i.eq.2)at=fn
        if(i.gt.2)print '(1p,5e15.7)',x,fn,(polv(i,j),j=1,3)
20 continue
c print '( ' Atilde= ' ',1p,e12.4, ' ' Btilde= ' ',e12.4, ' ' B/A = ' ',
c & e12.4)',at,bt,bt/at
    phi1=phi0+2*lam1*lsth
    z=cplx(ga/2,lam1)
    z=clngam(z)
    phi1=phi1-2*aimag(z)-atan(2*lam1/ga)
    phi1=phi1+(1-sing)*pi/2
    phi1=phi1/(2*pi)
    iphi1=phi1
    phi1=phi1-iphil
    if(phi1.gt.half)phi1=phi1-1
    if(phi1.lt.-half)phi1=phi1+1
    phi1=phi1*2*pi
    print '( ' check: ' ',15x,3f10.5,1p,2e10.3)',
&          shalf*sqrt(ga**2+4*lam1**2)*half,
&          phi1,half*log(bt/at)/pi,at,bt
    endif
c estimate of B/A = exp(2pi lambda)
    ba(min(itba,4))=exp(2*pi*lam1)
    if(itba.lt.4)print '(i3,f15.8)',itba,at21(itba)
    itba=itba+1
    recomp=0
    if(itba.eq.4)recomp=1
c recompute the homographic relation if new extreme data have been found
    if(itba.gt.4)then
        if(at21(4).lt.at21(2))then
            recomp=1
            at21(2)=at21(4)
            ba(2)=ba(4)
        endif
        if(at21(4).gt.at21(3))then

```

```

        recomp=1
        at21(3)=at21(4)
        ba(3)=ba(4)
    endif
endif
if(recomp.eq.1)then
    xx=(at21(1)-at21(2))/(1/ba(1)-1/ba(2))
    yy=(at21(1)-at21(3))/(1/ba(1)-1/ba(3))
    c3=(1/ba(3)-1/ba(2))/(xx/yy-1)-1/ba(2)
    c2=(at21(1)-at21(2))*(1/ba(1)+c3)*(1/ba(2)+c3)/(1/ba(2)-1/ba(1))
    c1=at21(1)-c2/(1/ba(1)+c3)
    print '(('',f15.6,'' A+''',f15.6,'' B)/(A+ ''',f15.6,'' B)''')',
&          c1,c1*c3+c2,c3
    print '(('' a1 tilde square min,max= ''',2f10.5)',c1,c1+c2/c3
endif
if(itba.eq.2)at21(2)=0
if(itba.eq.3)at21(3)=1
goto 11
end
end
/u18/grpanma/magnus/jacobi@ux12 [3] #a.out

```

```

    period of display of coeff., number of abs., extr. abs.?
50000 0 0 0
input alpha,beta,gamma,xtilde0,nmax (stop if nmax<=0)
0.1 -0.66 -0.9 0.8 500000
    0.10000000 -0.66000000 -0.90000000 0.80000000 500000
1 0.64000000 2 0.48642906 3 0.66612028
( 0.666359 A+ 0.682666 B)/(A+ 1.404457 B)
a1 tilde square min,max= 0.66636 0.48607
? lambda ( stop if > 999) 1
n atilde2(n) K phase lambda lambda
1 0.4863105269
2 0.1072928956 0.56893 0.63299 -0.23606 -0.11591
50000 0.2500104918 0.65794 2.04846 1.00002 1.00003
100000 0.2499993374 0.65795 2.04921 0.99999 1.00000
150000 0.2500017530 0.65794 2.04881 1.00001 1.00001
200000 0.2499986110 0.65795 2.04911 1.00000 1.00000
250000 0.2500025306 0.65795 2.04887 1.00001 1.00001
300000 0.2499980788 0.65795 2.04903 1.00000 1.00000
350000 0.2500018867 0.65796 2.04897 1.00000 1.00001
400000 0.2499994423 0.65795 2.04895 1.00000 1.00001
450000 0.2499999531 0.65795 2.04904 1.00000 1.00000
500000 0.2500011150 0.65795 2.04892 1.00001 1.00001
500001 0.2500009765 0.65795 2.04898 1.00000 1.00001
check: 0.65795 2.04898 1.00000 6.315E-05 3.382E-02
( 0.666359 A+ 0.682666 B)/(A+ 1.404457 B)

```

```

a1 tilde square min,max=    0.66636    0.48607
? lambda ( stop if > 999)    -2
  n    atilde2(n)          K      phase    lambda  lambda
    1  0.6663578695
    2  0.0057517038    0.88852  1.15619  -0.61116  -0.49518
 50000 0.2500181130    1.22984 -2.15364  -1.99944  -1.99940
100000 0.2500124101    1.22976 -2.15333  -1.99945  -1.99943
150000 0.2500050485    1.22973 -2.15338  -1.99945  -1.99944
200000 0.2499972423    1.22969 -2.15318  -1.99946  -1.99946
250000 0.2499956982    1.22970 -2.15294  -1.99947  -1.99947
300000 0.2500015654    1.22970 -2.15304  -1.99947  -1.99946
350000 0.2500034533    1.22971 -2.15326  -1.99946  -1.99945
400000 0.2499985243    1.22969 -2.15320  -1.99946  -1.99946
450000 0.2499982891    1.22970 -2.15304  -1.99947  -1.99946
500000 0.2500023390    1.22969 -2.15318  -1.99946  -1.99946
500001 0.2499993553    1.22970 -2.15306  -1.99947  -1.99946
check:                1.22969 -2.15322  -1.99946 4.756E-02 1.664E-07
(    0.666359 A+      0.682666 B)/(A+      1.404457 B)
a1 tilde square min,max=    0.66636    0.48607
? lambda ( stop if > 999)     3
  n    atilde2(n)          K      phase    lambda  lambda
    1  0.4860711314
    2  0.1073891907    0.56857  0.63221  -0.23556  -0.11549
 50000 0.2499658427    1.90240  1.71797  3.13855  3.13861
100000 0.2499817770    1.90236  1.71765  3.13858  3.13861
150000 0.2499992879    1.90235  1.71679  3.13862  3.13864
200000 0.2500085823    1.90243  1.71702  3.13861  3.13864
250000 0.2499927104    1.90240  1.71719  3.13861  3.13862
300000 0.2500063733    1.90242  1.71677  3.13862  3.13864
350000 0.2499951650    1.90239  1.71715  3.13861  3.13862
400000 0.2500046810    1.90242  1.71677  3.13862  3.13864
450000 0.2499960453    1.90241  1.71706  3.13861  3.13862
500000 0.2500039648    1.90243  1.71682  3.13862  3.13863
500001 0.2500009249    1.90243  1.71704  3.13862  3.13863
check:                1.90243  1.71683  3.13864 9.229E-11 3.386E-02
(    0.666359 A+      0.682666 B)/(A+      1.404457 B)
a1 tilde square min,max=    0.66636    0.48607
? lambda ( stop if > 999)    -4
  n    atilde2(n)          K      phase    lambda  lambda
    1  0.6663587554
    2  0.0057510998    0.88853  1.15619  -0.61116  -0.49518
 50000 0.2500178596    2.17836 -0.90688  -3.60233  -3.60219
100000 0.2499851215    2.17817 -0.90842  -3.60225  -3.60221
150000 0.2499982856    2.17817 -0.90898  -3.60222  -3.60219
200000 0.2499966056    2.17815 -0.90894  -3.60222  -3.60220
250000 0.2499928901    2.17814 -0.90874  -3.60223  -3.60221

```

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300000 0.2499936079 2.17816 -0.90855 -3.60224 -3.60222
350000 0.2500002389 2.17818 -0.90867 -3.60223 -3.60221
400000 0.2500054650 2.17815 -0.90898 -3.60222 -3.60220
450000 0.2500020824 2.17816 -0.90913 -3.60221 -3.60220
500000 0.2499962590 2.17814 -0.90891 -3.60222 -3.60221
500001 0.2500006090 2.17815 -0.90911 -3.60221 -3.60220
check: 2.17813 -0.90913 -3.60221 4.756E-02 7.042E-12
( 0.666359 A+ 0.682666 B)/(A+ 1.404457 B)
a1 tilde square min,max= 0.66636 0.48607
? lambda ( stop if > 999) 5
  n   atilde2(n)   K   phase   lambda   lambda
    1 0.4860711310
    2 0.1073891909 0.56857 0.63221 -0.23556 -0.11549
   26 -0.1229105534 9.79980 -0.13175 4.66254 4.89361
atildesquare <0, try a less extreme lambda
? lambda ( stop if > 999) 4.5
  n   atilde2(n)   K   phase   lambda   lambda
    1 0.4860711310
    2 0.1073891909 0.56857 0.63221 -0.23556 -0.11549
 50000 0.2500552560 2.75176 -0.28737 4.56431 4.56451
100000 0.2499780806 2.75188 -0.28663 4.56429 4.56436
150000 0.2499892472 2.75186 -0.28715 4.56432 4.56437
200000 0.2500094861 2.75176 -0.28772 4.56435 4.56439
250000 0.2500050796 2.75188 -0.28704 4.56432 4.56437
300000 0.2499915823 2.75187 -0.28722 4.56433 4.56435
350000 0.2500079170 2.75187 -0.28764 4.56435 4.56438
400000 0.2499943758 2.75181 -0.28714 4.56433 4.56435
450000 0.2500045768 2.75184 -0.28771 4.56435 4.56437
500000 0.2499965741 2.75183 -0.28722 4.56433 4.56435
500001 0.2499948277 2.75188 -0.28741 4.56434 4.56436
check: 2.75188 -0.28769 4.56436 1.188E-14 3.386E-02
( 0.666359 A+ 0.682666 B)/(A+ 1.404457 B)
a1 tilde square min,max= 0.66636 0.48607
? lambda ( stop if > 999) 5
  n   atilde2(n)   K   phase   lambda   lambda
    1 0.4860711310
    2 0.1073891909 0.56857 0.63221 -0.23556 -0.11549
 50000 0.2499526977 3.01064 -1.08241 4.99743 4.99760
100000 0.2500207551 3.01066 -1.08330 4.99749 4.99762
150000 0.2500144656 3.01067 -1.08379 4.99752 4.99761
200000 0.2499949927 3.01057 -1.08351 4.99751 4.99757
250000 0.2499902230 3.01067 -1.08422 4.99754 4.99758
300000 0.2500095958 3.01063 -1.08471 4.99757 4.99760
350000 0.2499959458 3.01060 -1.08402 4.99754 4.99757
400000 0.2500006139 3.01058 -1.08470 4.99757 4.99759
450000 0.2500018568 3.01067 -1.08429 4.99755 4.99758

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```
500000 0.2499976049 3.01065 -1.08462 4.99757 4.99758
500001 0.2500043587 3.01063 -1.08476 4.99757 4.99759
check:          3.01067 -1.08480 4.99756 7.809E-16 3.386E-02
( 0.666359 A+ 0.682666 B)/(A+ 1.404457 B)
a1 tilde square min,max= 0.66636 0.48607
? lambda ( stop if > 999) 9999
input alpha,beta,gamma,xtilde0,nmax (stop if nmax<=0)
0 0 0 0 0 0
STOP:
/u18/grpanma/magnus/jacobi@ux12 [4] #exit
exit
```

```
script done on Wed Mar 2 18:01:22 1994
```