# Stationary induction, reduction to the non-commutative Nevo–Zimmer theorem for *G*

#### Abstract

These are the notes for a presentation of about 1 hour in the joint seminar with Louvain-la-Neuve. This is the fifth lecture in the series.

# 1 Recap

The setting is as in the previous lectures. G is a connected simple Lie group with trivial center and rank  $\mathrm{rk}_{\mathbb{R}}(G) \geq 2$ . Then we can choose a maximal compact subgroup  $K \subset G$  and a maximal cocompact amenable subgroup  $P \subset G$  (minimal parabolic subgroup) such that G = KP. In this case, we denote  $\nu_P \in \mathrm{Prob}(G/P)$  the unique K-invariant (Borel) probability measure.

When  $\mu$  is a K-invariant measure on G, then the Poisson boundary of  $(G, \mu)$  is given by  $(G/P, \nu_P)$ . We can use uniqueness to see that  $\nu_P$  is  $\mu$ -stationary. Indeed, the measure  $\mu * \nu_P$  is K-invariant and therefore it must equal  $\nu_P$ .

We recall a few concepts from the previous lecture in the following definition.

**Definition 1.1.** Let  $\Lambda$  be a countable discrete group. Then  $PD_1(\Lambda)$  denotes the convex compact space of all the positive definite functions  $\varphi \colon \Lambda \to \mathbb{C}$  such that  $\varphi(e) = 1$ .  $\Lambda$  acts on  $PD_1(\Lambda)$  by conjugation, and we write  $Char(\Lambda) = PD_1(\Lambda)^{\Lambda}$ , so  $\varphi \in Char(\Lambda)$  if  $\varphi(\lambda x \lambda^{-1}) = \varphi(x)$  for all  $x, \lambda \in \Lambda$ .

If H is an lcsc group and  $\mu$  is a probability measure on H, then a bounded function  $f \in L^{\infty}(H,\mu)$  is called harmonic if

$$(\mu * f)(g) = \int_{H} f(gh)d\mu(h) = f(g)$$

for every  $g \in H$ . We denote by  $\operatorname{Har}^{\infty}(H, \mu)$  the space of bounded harmonic functions.

If  $H \curvearrowright M$  is an action on a von Neumann algebra, a state  $\psi$  on M is called  $\mu$ -stationary if  $(\mu * \psi)(x) = \int_H (h \cdot \psi)(x) d\mu(h) = \psi(x)$  for every  $x \in M$ . In that case we call  $(M, \psi)$  an  $(H, \mu)$ -von Neumann algebra. When M is abelian, this corresponds precisely to an  $(H, \mu)$ -space  $(B, \nu)$ .

**Lemma 1.2.** Suppose that  $(M, \psi)$  is an  $(H, \mu)$ -von Neumann algebra. Then we get a well-defined (normal ucp) H-equivariant map

$$\theta \colon M \to \operatorname{Har}^{\infty}(H, \mu) \colon \quad \theta_x(h) = (h \cdot \psi)(x).$$
 (1.1)

In particular Lemma 1.2 above applies to abelian von Neumann algebras, i.e. measure spaces  $(B, \nu)$  such that  $\nu$  is  $\mu$ -stationary. Recall that the *Poisson boundary*  $(B, \nu)$  is the  $(H, \mu)$  space such that the map (1.1) is a bijection.

*Proof.* It is straightforward to check that  $\theta$  is H-equivariant. Given  $x \in M$  we check that

 $\theta_x \in L^{\infty}(H,\mu)$  is harmonic. Indeed, using that  $\mu * \psi = \psi$ , we get that

$$(\mu * \theta_x)(g) = \int_H \theta_x(gh) d\mu(h)$$

$$= \int_H (gh \cdot \psi)(x) d\mu_h$$

$$= \int_H (h \cdot \psi)(g^{-1} \cdot x) d\mu(h)$$

$$= (\mu * \psi)(g^{-1} \cdot x) = \psi(g^{-1} \cdot x)$$

$$= (g \cdot \psi)(x) = \theta_x(g).$$

**Definition 1.3.** A probability measure  $\mu_0$  on  $\Gamma$  is called a *Furstenberg measure* if the following three conditions hold.

- 1.  $\mu_0(\gamma) > 0$  for every  $\gamma \in \Gamma$ .
- 2. The Poisson boundary of  $(\Gamma, \mu_0)$  equals  $(G/P, \nu_P)$ .

In particular we have that  $\mu_0 * \nu_P = \nu_P$ .

Let  $\Lambda$  be a countable discrete group. If  $\mu$  is a probability measure on  $\Lambda$  we say that  $\varphi \in PD_1(\Lambda)$  is a  $\mu$ -character if  $(\mu * \varphi) = \varphi$ , i.e.

$$\sum_{\lambda \in \Lambda} \mu(\lambda) \varphi(\lambda^{-1} x \lambda) = \varphi(x),$$

for every  $x \in \Lambda$ . Clearly, and character on  $\Lambda$  is a  $\mu$ -character on  $\Lambda$ .

**Theorem 1.4** (Theorem C). Let  $\Gamma \subset G$  be a lattice and let  $\mu_0$  be a Furstenberg measure on  $\Gamma$ . Then the following hold.

- 1. Any  $\mu_0$ -character on  $\Gamma$  is a genuine character.
- 2. Any extreme point  $\varphi \in \operatorname{Char}(\Gamma)$  either gives rise to a finite dimensional representation or  $\varphi = \delta_e$ .

**Theorem 1.5** (Theorem D). Let  $\Gamma \subset G$  be a lattice and let  $\mu_0$  be a Furstenberg measure on  $\Gamma$ . Let  $(M, \psi)$  be an ergodic  $(\Gamma, \mu_0)$ -von Neumann algebra. Then precisely one of the following holds.

- 1. Either  $\psi$  is  $\Gamma$ -invariant.
- 2. Or there exists a proper parabolic subgroup  $P \subset Q \subset G$  with  $Q \neq G$  and a state preserving  $\Gamma$ -equivariant normal unital \*-embedding

$$L^{\infty}(G/Q, \nu_Q) \to (M, \psi),$$

where  $\nu_Q$  denotes the unique K-invariant Borel probability measure on G/Q.

Our first goal is prove Theorem C using Theorem D!

**Remark 1.6.** We will apply Theorem D to the more specific situation where  $\pi \colon \Gamma \to \mathcal{U}(M)$  is a unitary representation and  $\sigma \colon \Gamma \curvearrowright M$  is given by  $\sigma_{\gamma}(x) = \operatorname{Ad}(\pi(\gamma))(x)$ . Then, if we happen to be in situation 2 of Theorem D, it follows from Lemma 2.2 that  $\psi \circ \pi = \delta_e$ . If time permits, I will try to give an idea of why this is true.

## 2 Deduce Theorem D from Theorem C

*Proof.* 1. The space  $\operatorname{Char}_{\mu_0}(\Gamma)$  of all  $\mu_0$ -characters on  $\Gamma$  is a nonempty compact convex, so by the Krein–Milman theorem, it suffices to prove the statement for an extreme point  $\varphi \in \operatorname{Char}_{\mu_0}(\Gamma)$ .

So let  $\varphi \in \operatorname{Char}_{\mu_0}(\Gamma)$  be an extreme point and let  $(\pi, \mathcal{H}, \xi)$  denote the corresponding GNS triple. Set  $M = \pi(\Gamma)'' \subset B(\mathcal{H})$  and let  $\psi$  be the state on M given by  $\psi(x) = \langle x\xi, \xi \rangle$ . Recall that using this notation we get that

$$\varphi = \psi \circ \pi. \tag{2.1}$$

Let  $\sigma \colon \Gamma \curvearrowright M$  be the action given by conjugation, i.e.  $\sigma_{\gamma}(x) = \pi(\gamma)x\pi(\gamma^{-1})$  for  $x \in M$ . Then we make the following observations that follow directly from (2.1).

- 1. The normal state  $\psi$  on M is  $\mu_0$ -stationary.
- 2. Our choses state  $\varphi$  is conjugation invariant iff  $\psi$  is  $\Gamma$ -invariant.

Basically, since  $\varphi \in \operatorname{Char}_{\mu_0}(\Gamma)$  is an extreme point, we get that the action  $\Gamma \curvearrowright M$  is ergodic. And so we are in the setting of Theorem D. If  $\varphi$  is not conjugation invariant, then  $\psi$  is not  $\Gamma$ -invariant and therefore we are in the second situation of Theorem D. By Remark 1.6 we get that  $\varphi = \psi \circ \pi = \delta_e$  is conjugation invariant anyway.

2. Suppose that  $\varphi \in \operatorname{Char}(\Gamma)$  is an extreme point. Denote by  $(\pi, \mathcal{H}, \xi)$  the GNS representation associated to  $\varphi$ . Again define  $M = \pi(\Gamma)'' \subset B(\mathcal{H})$  together with the faithful normal state (!)  $\psi(x) = \langle x\xi, \xi \rangle$ . Then M is a factor, because  $\varphi$  is an extreme point. Let  $J \colon \mathcal{H} \to \mathcal{H}$  be the conjugate unitary given by  $J(x\xi) = x^*\xi$  so that JMJ = M'. Define actions

$$\alpha \colon \Gamma \curvearrowright L^{\infty}(G/P) \ \overline{\otimes} \ B(\mathcal{H}) \colon \quad \alpha_{\gamma} = \sigma_{\gamma} \otimes \operatorname{Ad}(J\pi(\gamma)J)$$
$$\beta \colon \Gamma \curvearrowright L^{\infty}(G/P) \ \overline{\otimes} \ B(\mathcal{H}) \colon \quad \beta_{\gamma} = \operatorname{Ad}(1 \otimes \pi(\gamma)).$$

The actions  $\alpha$  and  $\beta$  commute, and thus we obtain a well-defined action

$$\beta \colon \Gamma \curvearrowright N = (L^{\infty}(G/P) \overline{\otimes} B(\mathcal{H}))^{\alpha}. \tag{2.2}$$

We can think of N as bounded measurable function  $f: G/P \to B(\mathcal{H})$  satisfying the equivariance property  $f(\gamma \cdot w) = \operatorname{Ad}(J\pi(\gamma)J)f(w)$  for every  $\gamma \in \Gamma$  and a.e.  $w \in G/P$ . Thus we can identify  $N \cong (L^{\infty}(G/\Gamma) \boxtimes B(\mathcal{H}))^P$ . Since P is amenable, we conclude that N is an amenable von Neumann algebra. Also note that  $1 \otimes M \subset N$ . Consider the normal state  $\Psi \in N_*$  given by

$$\Psi \colon N \to \mathbb{C} : f \mapsto \int_{G/P} \langle f(w)\xi, \xi \rangle d\nu_P(w).$$

Then  $\Psi$  is  $\mu_0$ -stationary, which follows from the stationarity of  $\nu_P$ . Furthermore, the action  $\beta \colon \Gamma \curvearrowright N$  is ergodic. Indeed, we may view an element  $f \in N^{\beta}$  as a  $\Gamma$ -equivariant function  $G/P \to M' = JMJ$ . Then we need to show that f is essentially constant to a scalar multiple of the identity. Define the measurable function

$$F: G/P \times G/P \to \mathbb{R}: \quad F(w_1, w_2) = ||Jf(w_1)J - Jf(w_2)J||_2,$$

where  $\|\cdot\|_2$  denotes the norm induced by the trace  $\psi$ . By  $\Gamma$ -equivariance of f we get that F is  $\Gamma$ -invariant. Since the Poisson boundary is doubly ergodic, we get that F is essentially constant. Then we conclude that f is essentially constant, say  $f(w) = x \in M'$  for a.e.  $w \in G/P$ . Again

by equivariance we get that x commutes with  $J\pi(\Gamma)J$ , so that  $x \in M \cap M' = \mathbb{C}1$ , which is what we needed to show.

Now we are again in the setting of Theorem D. If  $\Psi$  is not  $\beta$ -invariant, we get that  $\Psi \circ (1 \otimes \pi) = \delta_e$  by Remark 1.6. Note that  $\Psi$  restricts to  $1 \otimes \psi$  on  $1 \otimes M$ . Then it follows that  $\psi \circ \pi = \delta_e$ , i.e.  $\psi$  is  $\varphi = \delta_e$ .

If  $\Psi$  is  $\Gamma$ -invariant, then we show that  $1 \otimes M = N$ , so that M is amenable and has property (T). Then it follows that M is finite dimensional. Take  $f \in N$ , then one can show that

$$\Psi(\beta_{\gamma}^{-1}(f)) \int_{G/P} \langle f(\gamma \cdot w)\xi, \xi \rangle d\nu_P(\omega).$$

Thus writing  $\theta$  for the isometric bijection

$$\theta \colon L^{\infty}(G/P, \nu_P) \to \operatorname{Har}^{\infty}(\Gamma, \mu_0) : \quad \theta_h(g) = \int_{G/P} h(g \cdot w) d\nu_P(w).$$

we get that  $\Psi(\beta_{\gamma}^{-1}(f)) = \theta_h(\gamma)$ , where  $h \in L^{\infty}(G/P)$  is given by  $h(w) = \langle f(w)\xi, \xi \rangle$ . By invariance of  $\Psi$  we conclude that h is essentially constant. Since  $1 \overline{\otimes} M \subset N$ , we deduce that for every  $f \in N$  and every  $a, b \in M$  we get that  $(1 \otimes b^*)f(1 \otimes a) \in N$  and so the map  $w \mapsto \langle f(w)a\xi, b\xi \rangle$  is essentially constant. By density of  $M\xi \subset \mathcal{H}$  we get that  $w \mapsto f(w)$  is essentially constant, and equal to some element  $x \in B(\mathcal{H})$  say. By the equivariance of f we get again that  $x \in M$ , and so we get the desired conclusion  $f \in 1 \otimes M$ .

### 3 On Lemma 2.2

Suppose that  $\Gamma \to \mathcal{U}(M)$  is a unitary representation, so that the action  $\sigma \colon \Gamma \curvearrowright M$  given by  $\sigma_{\gamma} = \operatorname{Ad}(\pi(\gamma))$  makes  $(M, \psi)$  into an ergodic  $(\Gamma, \mu_0)$ -von Neumann algebra. Assume moreover there exists a proper parabolic subgroup  $P \subset Q \subset P$  and a state preserving unital \*-embedding  $L^{\infty}(G/Q, \nu_Q) \to M$ .

Let  $A \subset M$  be the separable unital  $C^*$ -algebra generated by  $\pi(\Gamma)$  and  $\theta(C(G/Q))$ . Then A is globally  $\Gamma$ -invariant and the state  $\psi|_A$  is  $\mu_0$ -stationary. Then we saw in the previous lecture that there exists an essentially unique map  $G/P \to \mathcal{S}(A) : w \mapsto \psi_w$  such that

$$\psi = \int_{G/P} \psi_w d\nu_P(w). \tag{3.1}$$

Denote  $p_Q: G/P \to G/Q$  the measure preserving factor map. The state  $\psi: \theta$  is  $\mu_0$ -stationary, and thus it must equal  $\nu_Q$ . Then we get that

$$\int_{G/P} \psi_w \circ \theta d\nu_P(w) = \psi \circ \theta = \delta_Q = \int_{G/P} \delta_{p_Q(w)} d\nu_P(w).$$

By uniqueness of boundary maps we get that  $\psi_w \circ \theta = \delta_{p_Q(w)}$  for a.e.  $w \in G/P$ . As  $\delta_{p_Q(w)}$  is multiplicative, we get that  $\theta(C(G/Q))$  lies in the multiplicative domain of  $\psi_w$  for a.e.  $w \in G/P$ .

Pick  $\gamma \neq e$ . Since the action  $\Gamma \curvearrowright G/Q$  is essentially free and since  $p_Q \colon G/P \to G/Q$  is pmp we get that  $\gamma \cdot p_Q(w) \neq p_Q(w)$  for a.e. G/P. Pick  $w \in G/P$  such that  $\gamma \cdot p_Q(w) \neq p_Q(w)$  and such that  $\theta(C(G/Q))$  is in the multiplicative domain of  $\psi_w$  and write  $y = \pi_Q(w)$ . Take  $f \in C(G/Q)$  such that f(y) = 1 and  $f(\gamma \cdot y) = 0$ . Then we get that

$$\psi_w(\pi(\gamma)) = f(y)\psi_w(\pi(\gamma) = \psi_w(\theta(f)\pi(\gamma)) = \psi_w(\pi(\gamma)\pi(\gamma^{-1})\theta(f)\pi(\gamma))$$
$$= \psi_w(\pi(\gamma)\theta(\gamma^{-1} \cdot (f))) = \psi_w(\pi(\gamma))f(\gamma \cdot y) = 0.$$

Since this holds for a.e.  $w \in G/P$  we conclude from (3.1) that  $\psi \circ \pi = \delta_e$ .