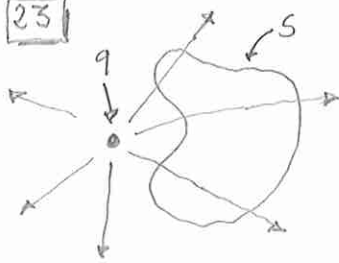


23



$$\oint \vec{E} \cdot d\vec{S} = \frac{\sum q_i}{\epsilon_0}$$

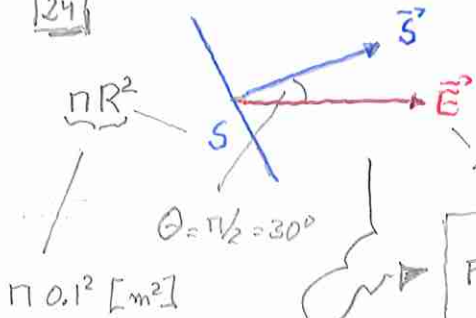
= 0

$$\sum q_i = 0$$

FLUX = 0
 \Rightarrow PAS DE CHARGES !

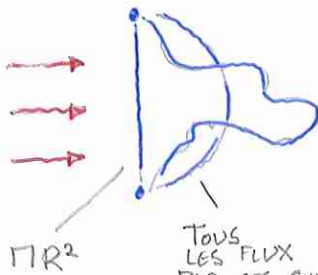
FLUX = 0
 $\nRightarrow \vec{E} = 0$
 HEHE! si $\vec{E} = 0 \Rightarrow$ FLUX 0 :-)

24



$$\text{FLUX} = S E \cos(\varphi) = \pi/2 \text{ [Nm}^2/\text{C]}$$

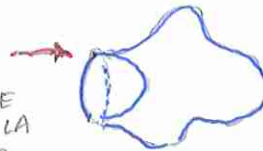
25



$$\text{FLUX} = \pi R^2 E$$

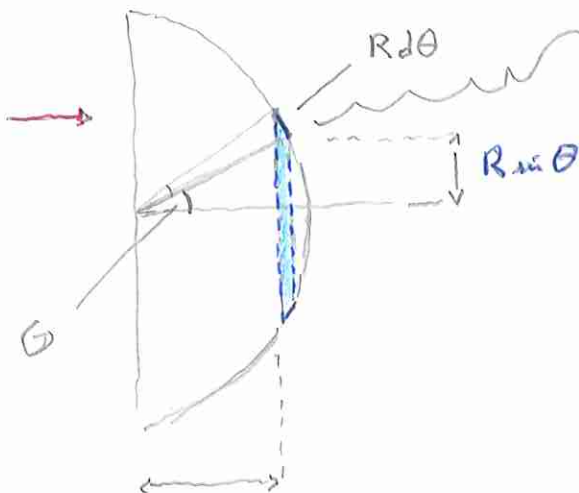
TOUS LES FLUX PAR CES SURFACES SONT IDENTIQUES

- LE CERCLE OU LA DEMI-SPHERE OU N'IMPORTE QUELLE SURFACE QUI A LA MEME OUVERTURE



YES SO EASY!

• ET SI TU N'ES PAS CONVAINCU ET QUE TU AIMES L'ALGÈBRE !

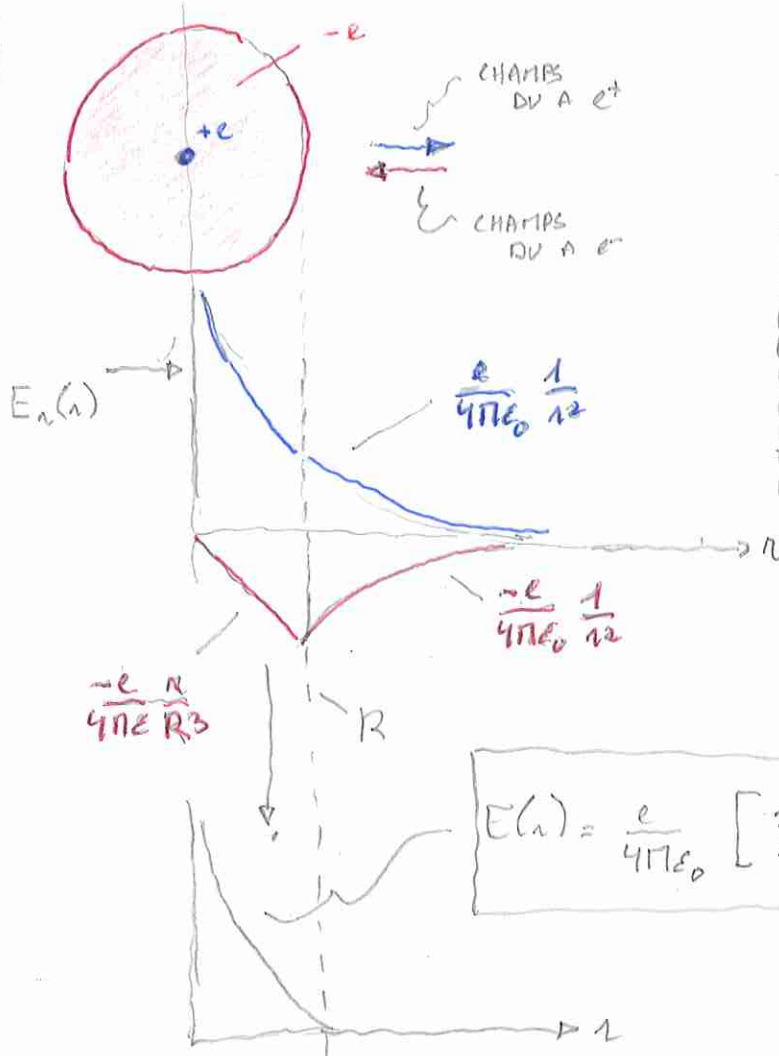


$$d\text{FLUX} = \underbrace{2\pi R \sin \theta R d\theta}_{\text{SURFACE COURONNE } dS \text{ :-)}} \cos(\theta) E$$

$$\text{FLUX} = 2\pi R^2 E \int_0^{\pi/2} \sin^2 \theta \cos \theta d\theta$$

$$= \pi R^2 E \left[\frac{\sin^3 \theta}{3} \right]_0^{\pi/2}$$

26



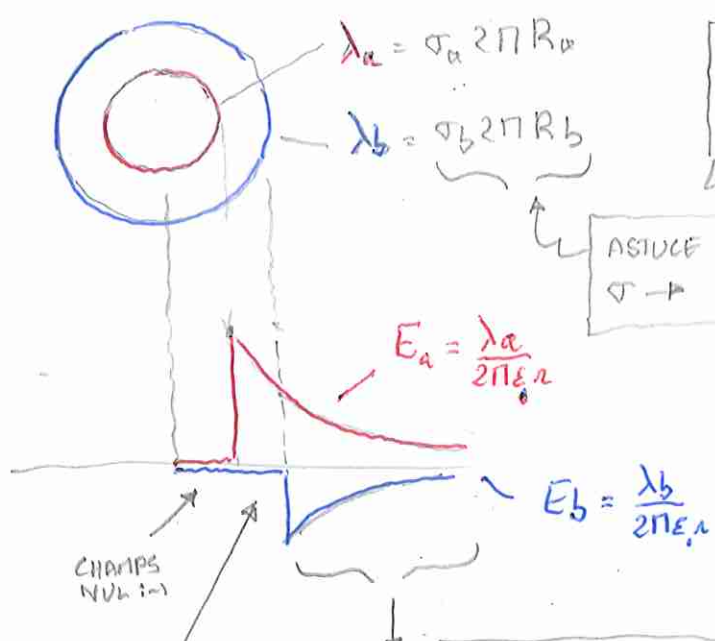
IL SUFFIT DE SUPERPOSER LES CHAMPS D'UNE CHARGE PONCTUELLE $+e$ ET D'UNE CHARGE $-e$ REPARTIE SUR LA SPHERE

APPLICATION CAS PARTICULIER #1

$$E(r) = \frac{e}{4\pi\epsilon_0} \left[\frac{1}{r^2} - \frac{r}{R^3} \right]$$

APPLICATION CAS PARTICULIER #2

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ON SUPERPOSE LES CHAMPS DE DEUX LIGNES CHARGES !

ASTUCE ! $\sigma \rightarrow \lambda$



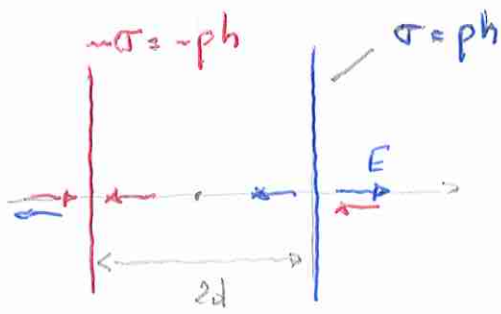
CHAMPS NUL A L'EXTERIEUR $\lambda_a = -\lambda_b$

$r \in [R_a, R_b]$

$$E(r) = \frac{1}{2\pi\epsilon_0} \frac{R_a \sigma_a}{r}$$

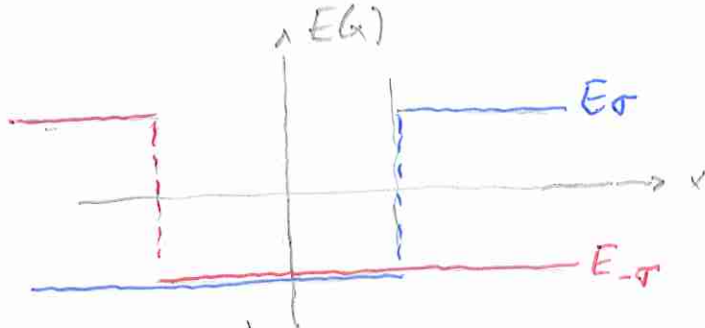
$$R_a \sigma_a = -\sigma_b R_b$$

28

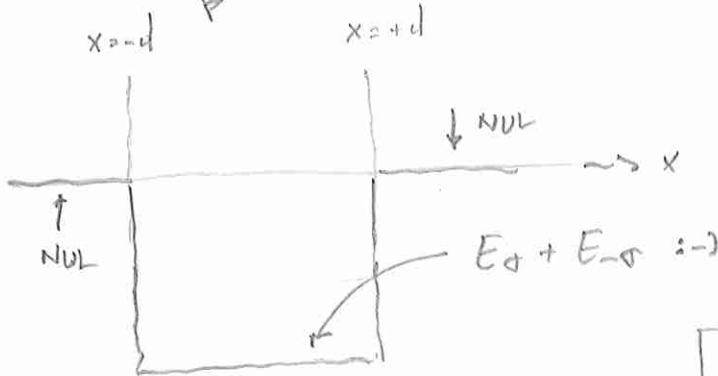


ON SUPERPOSE
LES CHAMPS DE DEUX
PLANS CHARGÉS :-)

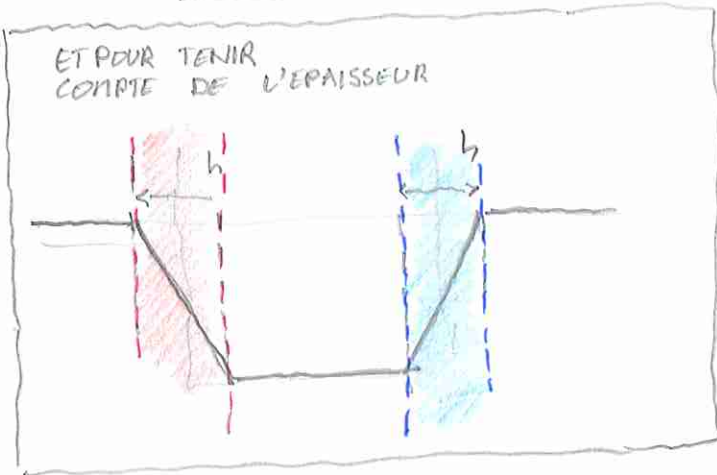
APPLICATION
CAS PARTICULIER #3



EN
ADDITIONNANT :-)



TOUS
LES EXERCICES
SONT
TOUJOURS
CONSTRUITS SUR
LE MEME
PRINCIPE !



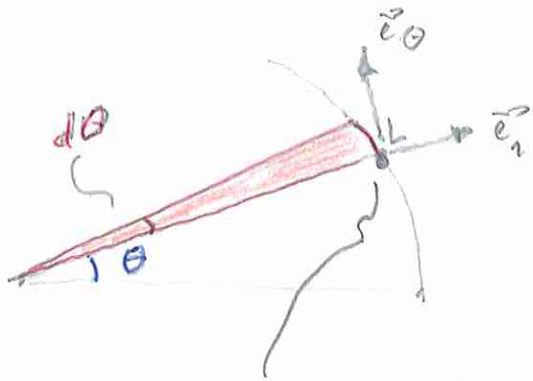
ASTUCE :-)
CONDUCTEUR
CHAMP
À L'INTÉRIEUR = 0 :-)
ISOLANT
VARIATION
LINÉAIRE
DU CHAMP :-)

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$$\oint_S \vec{g} \cdot d\vec{S} = -4\pi G \sum m_i$$

↑
L'ATTRACTION
EST UNE FORCE
ATTRACTIVE :-)

LA FORCE ÉLECTRIQUE
ENTRE 2 CHARGES IDENTIQUES
EST REPULSIVE



$$\vec{e}_r = \begin{bmatrix} \cos \theta \\ \sin \theta \end{bmatrix}$$

$$\vec{e}_\theta = \begin{bmatrix} -\sin \theta \\ \cos \theta \end{bmatrix}$$

$$\vec{r}(t) = \begin{bmatrix} r(t) \cos(\theta(t)) \\ r(t) \sin(\theta(t)) \end{bmatrix}_{xy} = x_x \vec{e}_x + x_y \vec{e}_y$$

CHAIN'S RULE!

$$\vec{v}(t) = \begin{bmatrix} r' \cos \theta - r \theta' \sin \theta \\ r' \sin \theta + r \theta' \cos \theta \end{bmatrix}_{xy}$$

$$\vec{a}(t) = \begin{bmatrix} r'' \cos \theta - 2r' \theta' \sin \theta - r \cos \theta (\theta')^2 - r \sin \theta \theta'' \\ r'' \sin \theta + 2r' \theta' \cos \theta - r \sin \theta (\theta')^2 + r \cos \theta \theta'' \end{bmatrix}_{xy}$$

$$= [r'' - r(\theta')^2] \vec{e}_r + [2r'\theta' + r\theta''] \vec{e}_\theta$$

oui :-)

COMME L'UNIQUE FORCE EST SELON LA DIRECTION RADIALE !

LA COMPOSANTE a_θ EST NULLE

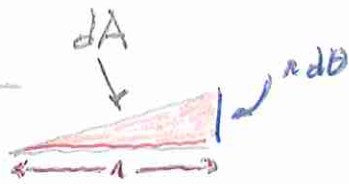
$$2 \frac{dr}{dt} \frac{d\theta}{dt} + r \frac{d^2\theta}{dt^2} = 0$$

$$2r \frac{dr}{dt} \frac{d\theta}{dt} + r^2 \frac{d^2\theta}{dt^2} = 0$$

$$\frac{d}{dt} \left[r^2 \frac{d\theta}{dt} \right] = 0$$

CE TERME EST CONSTANT = k

CALCULONS L'AIRE :-)



$$A = \int_{\theta}^{\theta + \Delta t} \frac{1}{2} r(\theta) r(\theta) d\theta$$

$$= \frac{1}{2} \int_{t_1}^{t_1 + \Delta t} r^2(t) \frac{d\theta(t)}{dt} dt$$

$$= \frac{1}{2} k \Delta t$$